

# Edgeworth Expansions

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## ◆ Outline

- Basic Idea
- Charlier Differential Series
- Edgeworth Expansions
- Applications
  - Numerical Examples
  - Bootstrapping Studentized Statistics
- Conclusions

Try to improve *CLT* by approx.  
 $F_n(t)$  by  $\sum_{j=0}^{\infty} \frac{A_j(t)}{n^{j/2}}$  such that

$$\left| F_n(t) - \sum_{j=0}^r \frac{A_j(t)}{n^{j/2}} \right| \leq \frac{C_r(t)}{n^{(r+1)/2}}$$

- $\lim_{r \rightarrow \infty} \sum_{j=0}^r \frac{A_j(t)}{n^{j/2}}$  may or may not exist for fixed  $n$ ;
- error is of the same order of magnitude as the first neglected term

## ♦ Basic Idea

Can try to improve the accuracy by using  
e.g. Edgeworth expansions;  
cf. Feller(1971), Ch. 16.

They are obtained by a Taylor expansion of the characteristic function of the statistic of interest around 0, i.e. at the center of the distribution, followed by a Fourier inversion.

This leads to expansions of the distribution in powers of  $n^{-1/2}$ , where the leading term is the normal density.

## ♦ Charlier Differential Series (1905)

|               |           |   |
|---------------|-----------|---|
| Distr.        | $H(x)$    | $G(x)$  |
| Charact. fct. | $\chi(u)$ | $\xi(u) = \int e^{iux} dG(x)$                           |
| Cumulants     | $\beta_r$ | $\gamma_r = (-i)^r \frac{d^r}{du^r} \log \xi(u) _{u=0}$ |

Let us first express  $H(x)$  through  $G(x)$  by means of their characteristic functions:

$$\chi(u) = \exp\left\{\sum_{r=1}^{\infty} (\beta_r - \gamma_r) \frac{(iu)^r}{r!}\right\} \xi(u)$$

Suppose all derivatives of  $G$  vanish at the extremes of the range of  $x$ . Then, by integration by parts

$(iu)^r \xi(u)$  is the Fourier transform of  $(-1)^r G^{(r)}(x)$

Then, by Fourier inversion:

$$H(x) = \exp\left\{\sum_{r=1}^{\infty} (\beta_r - \gamma_r) \frac{(-D)^r}{r!}\right\} G(x)$$

D: differential operator which denotes differentiation with respect to  $x$ .

- $r = 1$ :

$$\begin{aligned}\xi(u) &= \int_{-\infty}^{+\infty} e^{iux} g(x) dx \\ &= \underbrace{\left[ \frac{1}{iu} e^{iux} g(x) \right]_{-\infty}^{+\infty}}_{= 0} - \frac{1}{iu} \int_{-\infty}^{+\infty} e^{iux} g'(x) dx\end{aligned}$$

Thus,

$$(iu)\xi(u) = \int_{-\infty}^{+\infty} e^{iux} (-g'(x)) dx$$

is the **Fourier inverse** of  $-g'(x)$ .

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$$\begin{aligned}y &= a_1 x + a_2 x^2 + \dots \\ e^y &= 1 + y + \frac{y^2}{2} + \frac{y^3}{3!} + \dots\end{aligned}$$

## ♦ Edgeworth Expansions

$$\begin{array}{llll}
 X_1, \dots, X_n \text{ iid} & E(X_i) = 0 & = \chi_1(X_i) & \\
 & \text{var}(X_i) = \sigma^2 & = \chi_2(X_i) & \\
 & r \geq 3 & \chi_r(X_i) & = \lambda_r \sigma^r
 \end{array}$$

$$F_n(t) = P[T_n < t] \ , \quad T_n = \frac{1}{n^{1/2}} \sum_{i=1}^n \frac{X_i}{\sigma}$$

Edgeworth uses Charlier differential series to approx.  $F_n(t)$  by  $\Phi(t)$  (cumul. normal) expanding  $\exp\{\dots\}$  and collecting terms according to powers of  $\frac{1}{n^{1/2}}$ .



$$\begin{array}{ll}
H(\cdot) = F_n(\cdot) & \\
\chi(\cdot) = \psi_n(\cdot) & \\
\beta_1 = \chi_1(T_n) = 0 & = \\
\beta_2 = \chi_2(T_n) = 1 & = \\
\beta_r = \chi_r(T_n) = \frac{\lambda_r}{n^{r/2-1}} & 
\end{array}
\begin{array}{l}
G(\cdot) = \Phi(\cdot) \\
\xi(u) = \exp\{-u^2/2\} \\
\gamma_1 = 0 \\
\gamma_2 = 1 \\
\gamma_r = 0 \quad r \geq 3
\end{array}$$

$$\begin{aligned}
\psi_n(u) &= \exp \left\{ \sum_{r=3}^{\infty} \frac{\lambda_r}{n^{r/2-1}} \frac{(iu)^r}{r!} \right\} \cdot e^{-u^2/2} \\
&= \exp \left\{ \frac{\lambda_3}{n^{1/2}} \frac{(iu)^3}{3!} + \frac{\lambda_4}{n} \frac{(iu)^4}{4!} + \frac{\lambda_5}{n^{3/2}} \frac{(iu)^5}{5!} \right. \\
&\quad \left. + \dots \right\} \cdot e^{-u^2/2}
\end{aligned}$$

Collecting terms of order  $\frac{1}{n^{1/2}}, \frac{1}{n}, \frac{1}{n^{3/2}}, \dots$ , we get:

$$\begin{aligned}
\psi_n(u) &= \left\{ 1 \right. \\
&\quad + \frac{1}{n^{1/2}} \lambda_3 \frac{(iu)^3}{6} \\
&\quad + \frac{1}{n} \left[ \frac{1}{2} \lambda_3^2 \frac{(iu)^6}{(3!)^2} + \lambda_4 \frac{(iu)^4}{4!} \right] \\
&\quad \left. + \frac{1}{n^{3/2}} \dots \right\} \cdot e^{-u^2/2}
\end{aligned}$$

By Fourier inversion, we obtain the Edgeworth expansion for the cumulative distribution:

$$\begin{aligned}
 F_n(t) &= \Phi(t) \\
 &- \frac{1}{n^{1/2}} \frac{\lambda_3}{6} \Phi^{(3)}(t) \\
 &+ \frac{1}{n} \left[ \frac{\lambda_4}{24} \Phi^{(4)}(t) + \frac{\lambda_3^2}{72} \Phi^{(6)}(t) \right] \\
 &+ \dots
 \end{aligned}$$

For the density:

$$\begin{aligned}
 f_n(t) &= \phi(t) \\
 &- \frac{1}{n^{1/2}} \frac{\lambda_3}{6} \phi^{(3)}(t) \\
 &+ \frac{1}{n} \left[ \frac{\lambda_4}{24} \phi^{(4)}(t) + \frac{\lambda_3^2}{72} \phi^{(6)}(t) \right] \\
 &+ \dots \\
 &= \phi(t) \left\{ 1 \right. \\
 &+ \frac{1}{n^{1/2}} \frac{\lambda_3}{6} H_3(t) \\
 &+ \frac{1}{n} \left[ \frac{\lambda_4}{24} H_4(t) + \frac{\lambda_3^2}{72} H_6(t) \right] \\
 &+ \dots \left. \right\}
 \end{aligned}$$

where

$$H_3(t) = t^3 - 3t$$

$$H_4(t) = t^4 - 6t^2 + 3$$

$$H_6(t) = t^6 - 15t^4 + 45t^2 - 15$$

are the Hermite polynomials.

## Remarks

- $\frac{1}{n^{1/2}}$  term ( $\lambda_3$ ) corrects for skewness and  $\frac{1}{n}$  term ( $\lambda_4$ ) for kurtosis.

- In particular:

$$f_n(0) = \frac{1}{\sqrt{2\pi}} \left\{ 1 + \frac{1}{n} \left[ \frac{1}{8} \lambda_4 - \frac{5}{24} \lambda_3^2 \right] + O(n^{-2}) \right\}$$

- If the first  $K$  moments exist, then as  $n \rightarrow \infty$ :

$$F_n(t) - \Phi(t) - \sum_{r=3}^K \frac{P_r(-\Phi(t))}{n^{r/2-1}} = o\left(\frac{1}{n^{K/2-1}}\right),$$

where  $P_r(\cdot)$  is a polynomial of degree  $3(r-2)$  with coefficients depending only on  $\lambda_3, \dots, \lambda_r$ .

- The Edgeworth approximation is not a distribution function:  
0 – 1 range and monotonicity can be violated in parts of one or both tails.  
Numerical instability.
- Fisher-Cornish expansion for quantiles.
- Edgeworth expansion with respect to (e.g.)  $\chi^2$  distribution instead of the normal  
→ Laguerre polynomials instead of Hermite polynomials.

## Some References

Wallace(1958)

Feller(1971)

Bickel (1974), *Ann. Stat*

Albers, Bickel, van Zwet (1976), *Ann. Stat*

For  $L$ – statistics:

Helmers (1979, 1980), *Ann. Stat.*

For  $U$ – statistics:

Callaert, Janssen, Veraverbeke (1980), *Ann. Stat.*

Bickel, Götze, van Zwet (1986), *Ann. Stat.*

Multivariate Edgeworth expansions:

Skovgaard (1986), *Int. Stat. Rev.*

## ♦ Applications

### Numerical Examples

See Chapter 2 in

Field, C. and Ronchetti, E. (1990)

*Small Sample Asymptotics*

IMS Lectures Notes



## Bootstrapping Studentized Statistics

We want to compare the bootstrap distribution of a statistic  $\hat{\theta}$  with its studentized version.

Edgeworth expansions for the original statistic:

$$P_F[\sqrt{n}(\hat{\theta} - \theta) \leq t] = \Phi\left(\frac{t}{\sigma}\right) + \frac{1}{\sqrt{n}}p\left(\frac{t}{\sigma}\right)\phi\left(\frac{t}{\sigma}\right) + O\left(\frac{1}{n}\right) \quad (1)$$

$$P_F[\sqrt{n}\left(\frac{\hat{\theta} - \theta}{\hat{\sigma}}\right) \leq t] = \Phi(t) + \frac{1}{\sqrt{n}}q(t)\phi(t) + O\left(\frac{1}{n}\right) \quad (2)$$

Edgeworth expansions for the bootstrap version:

$$P_{\hat{F}}[\sqrt{n}(\hat{\theta}^* - \hat{\theta}) \leq t] = \Phi\left(\frac{t}{\hat{\sigma}}\right) + \frac{1}{\sqrt{n}}p\left(\frac{t}{\hat{\sigma}}\right)\phi\left(\frac{t}{\hat{\sigma}}\right) + O_p\left(\frac{1}{n}\right) \quad (3)$$

$$P_{\hat{F}}[\sqrt{n}\left(\frac{\hat{\theta}^* - \hat{\theta}}{\hat{\sigma}^*}\right) \leq t] = \Phi(t) + \frac{1}{\sqrt{n}}\hat{q}(t)\phi(t) + O_p\left(\frac{1}{n}\right) \quad (4)$$

Comparison:

$$\begin{aligned}(3) - (1) &= O_p\left(\Phi\left(\frac{t}{\hat{\sigma}}\right) - \Phi\left(\frac{t}{\sigma}\right)\right) \\&= O_p(\hat{\sigma} - \sigma) \\&= O_p\left(\frac{1}{\sqrt{n}}\right) \\(4) - (2) &= O_p\left(\frac{1}{\sqrt{n}}(\hat{q}(t) - q(t))\right) \\&= O_p\left(\frac{1}{\sqrt{n}}(\hat{\sigma}^* - \hat{\sigma})\right) \\&= O_p\left(\frac{1}{n}\right)\end{aligned}$$

Bootstrapping the studentized statistic captures the  $\frac{1}{\sqrt{n}}$  term of the Edgeworth expansion and provides second order correctness.

Hall(1986, 1988)

## ◆ Conclusions

- Useful theoretical tool to study the higher-order properties of estimators and test statistics.
- By construction Edgeworth expansions provide in general a good approximation in the center of the density, but they can be inaccurate in the tails, where they can even become negative.
- Absolute error.